

# Lane-Emden problem: algorithm and symmetries

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Let  $\Omega \subseteq \mathbb{R}^N$  be open bounded and  $2 < p < 2^*$ , where  $2^* = \frac{2N}{N-2}$  for  $N > 2$  and  $2^* = +\infty$  if  $N = 2$ .

## Lane-Emden problem:

$$\begin{cases} -\Delta u(x) = |u(x)|^{p-2}u(x), & x \in \Omega, \\ u(x) = 0, & x \in \partial\Omega. \end{cases}$$

Solutions are critical points of **energy** functional

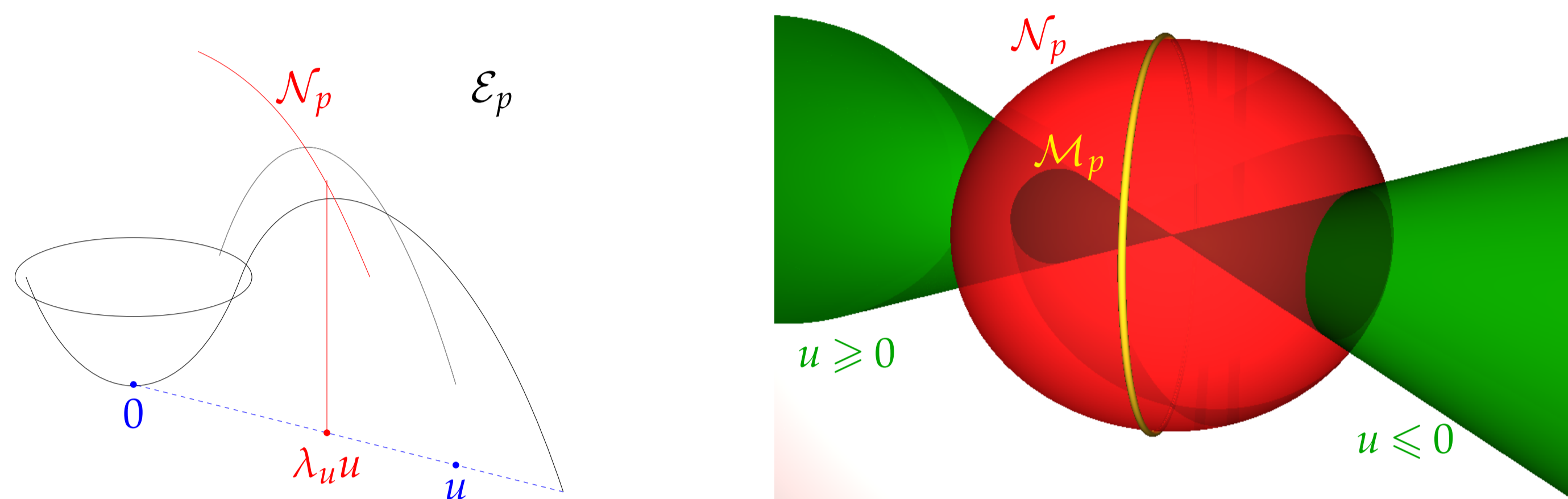
■  $\mathcal{E}_p : H_0^1(\Omega) \rightarrow \mathbb{R} : u \mapsto \frac{1}{2} \int_{\Omega} |\nabla u|^2 - \frac{1}{p} \int_{\Omega} |u|^p$ , where  $H_0^1(\Omega) := \overline{C_0^\infty(\Omega)}$  for  $\|u\|^2 = \int_{\Omega} |\nabla u|^2$ .

## Existence of GS and LENS:

**Theorem 1:** (Mountain-Pass theorem, [1] and [6])

There exist an one-signed solution and a sign-changing solution.

**Sketch:** Minima of  $\mathcal{E}_p$  on  $\mathcal{N}_p := \{u \neq 0 : \mathcal{E}'_p(u)(u) = 0\}$  and  $\mathcal{M}_p := \{u : u^\pm \in \mathcal{N}_p\}$ .



## Mountain-Pass Algorithm: ([7] and [14])

1. Let  $u \in H_0^1$  and  $n \leftarrow 0$ ;
2. Compute  $u_n \leftarrow P(u)$ ;
3. Compute  $g \leftarrow \nabla \mathcal{E}(u_n)$ ;
4. If  $\|\nabla \mathcal{E}(u_n)\| \leq \varepsilon$  stop; else  $v \leftarrow P(u_n - g)$ ;
5. If  $\mathcal{E}(v) < \mathcal{E}(u_n)$ ,  $u_{n+1} \leftarrow v$ ,  $n \leftarrow n + 1$  and go to step 3; else  $g \leftarrow \frac{g}{2}$  and go to step 4;

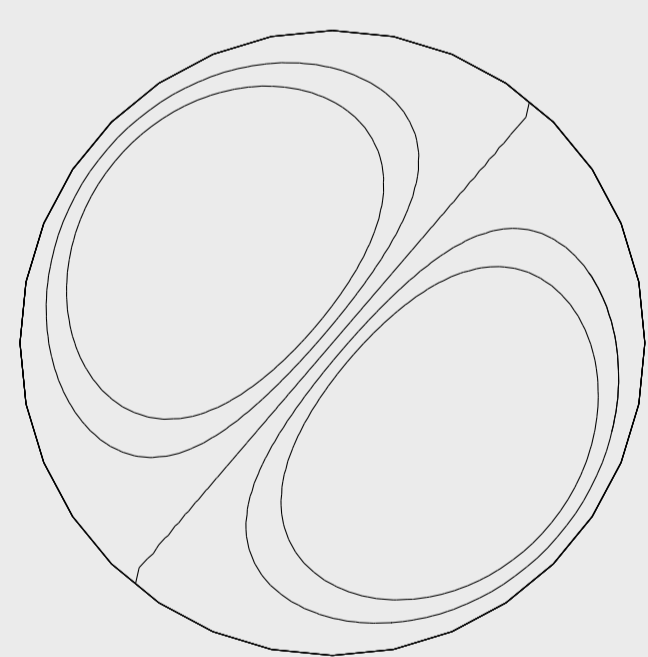
where the projection  $P$  equals  $P_N : H_0^1 \rightarrow \mathcal{N}_p : u \mapsto t^*u$  s.t.  $\mathcal{E}_p(t^*u) = \max_{t>0} \mathcal{E}_p(tu)$  to approach one-signed solution and  $P$  equals  $P_M : H_0^1 \rightarrow \mathcal{M}_p : u \mapsto P_N(u^+) + P_N(u^-)$  else.

**Remark:** A proof of convergence exists only for  $P = P_N$ .

## Symmetries: theoretical results

- On convex domains, **GS** respect symmetries of  $\Omega$  [8]. On non-convex domains, it is **not** always the case [5].
- On balls, for  $p$  small, **LENS** are even with respect to  $N - 1$  directions and odd with respect to the last one [3,10]. For  $p$  large, there only exists proof for even directions [2].
- On rectangles, for  $p$  small, **LENS** are even and odd with respect to a median. For  $p$  large, it is **not** always the case. We obtain a **symmetry breaking** [3].
- On squares, for  $p$  small, **LENS** are odd with respect to the center [3]. Numerically, they seem to be odd with respect to a diagonal.

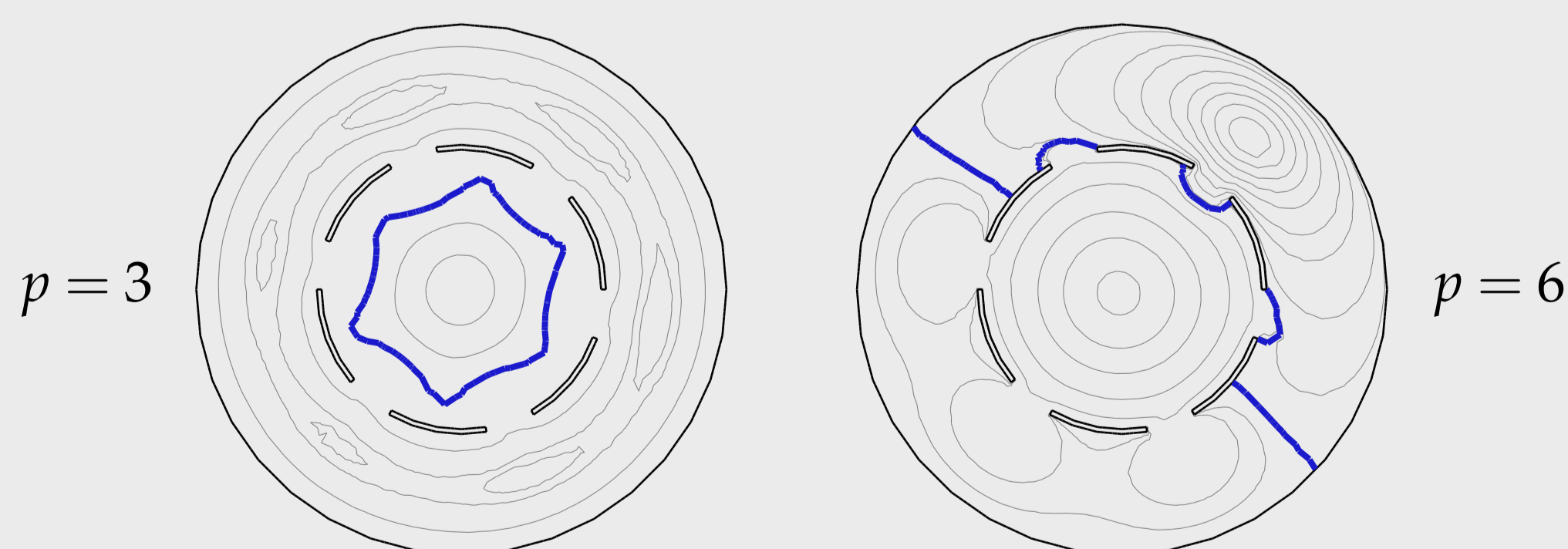
## Level curves and nodal line of LENS



On a ball, for  $p$  small, the nodal line is a diameter [3,10].

When the domain is convex and  $p$  close to 2, the nodal line **touches the boundary** [12]. It is conjectured to still be true for simply connected domains.

For not connected domains, the nodal line on the domain  $\Omega$  below does *not* touch  $\partial\Omega$  for  $p$  small [12]. Numerically, it netherless seems touch  $\partial\Omega$  for larger  $p$ .



## References

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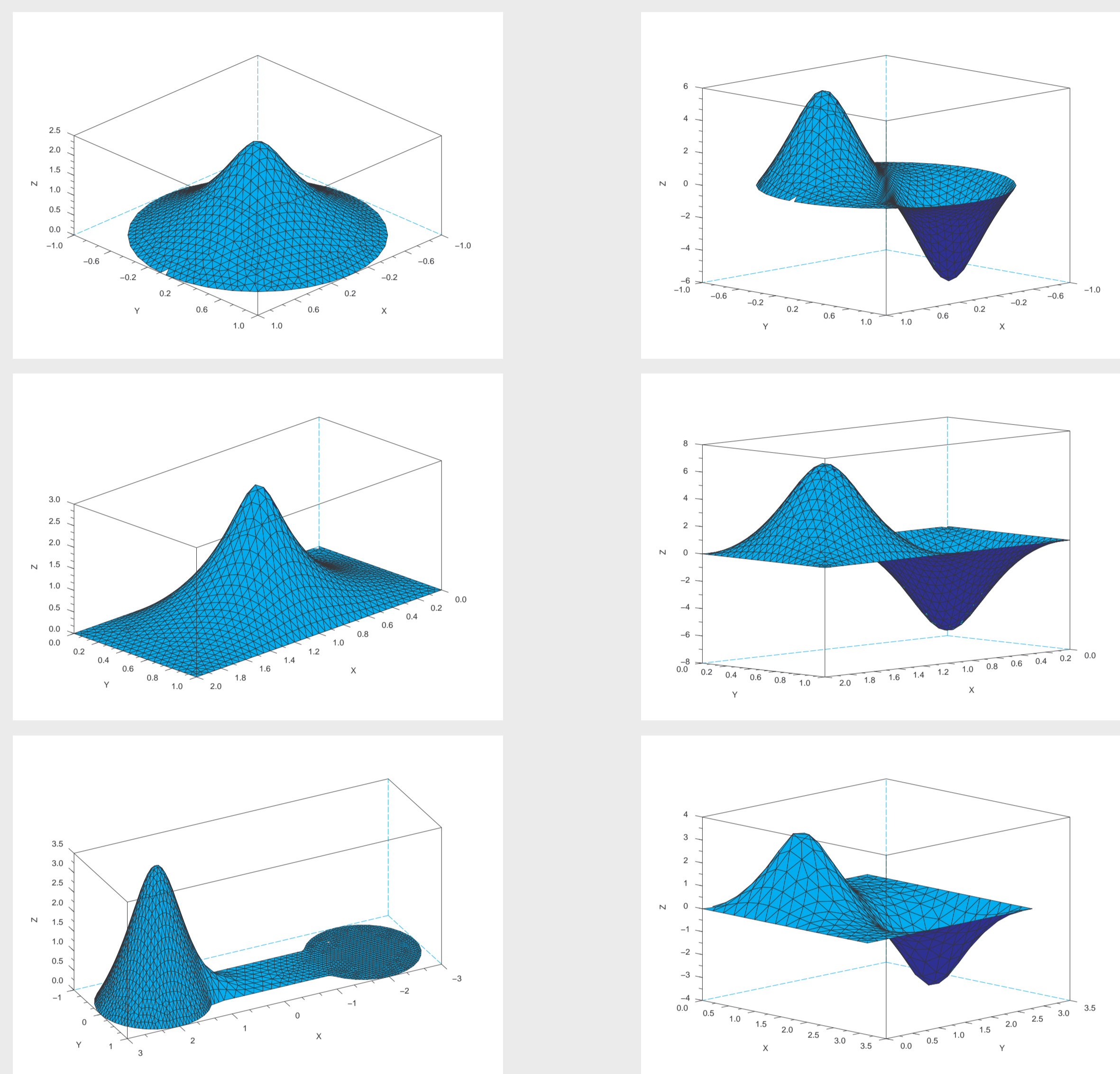
## Definitions

- **Ground state solution (GS):** non-zero one-signed solution with minimum energy
- **Least energy nodal solution (LENS):** sign-changing solution with minimum energy
- $u^+ := \max(u, 0)$  and  $u^- := \min(u, 0)$
- **Nodal line of  $u$  :**  $\text{adh}\{x \in \Omega : u(x) = 0\}$

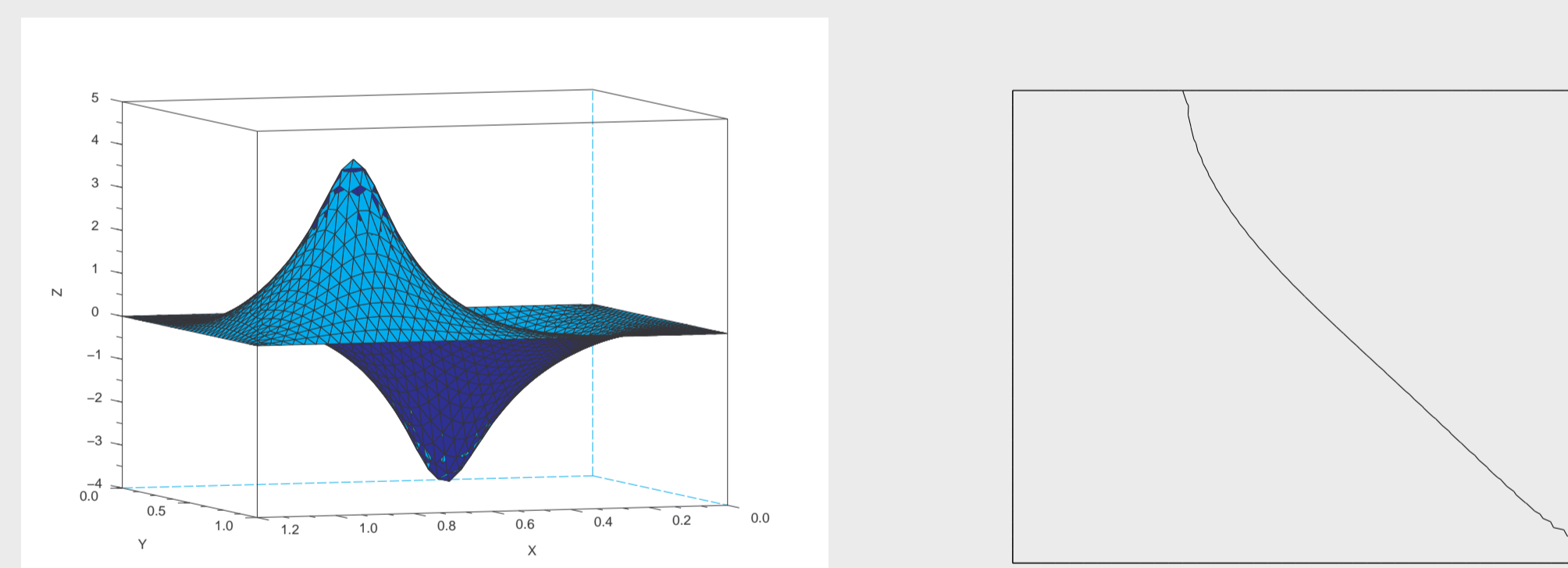
## Main objectives

- Existence of **GS** and **LENS**
- **Algorithm** to approach solutions
- Theoretical study of **symmetries** of **GS** and **LENS**

## Numerical solutions for $-\Delta u = u^3$



## Numerical solution and nodal line for $-\Delta u = u^5$



## Generalizations

- Previous algorithm can be generalized to approach **GS** of problem

$$\begin{cases} -\Delta u(x) + \mathbf{V}(x)u(x) = |u(x)|^{p-2}u(x), & x \in \Omega, \\ u(x) = 0, & x \in \partial\Omega, \end{cases}$$

when 0 does not belong to the spectrum  $\sigma(-\Delta + V)$  [11,13].

**Ideas:** Solutions are critical points of energy functional

$$\mathcal{E}_p : H_0^1(\Omega) \rightarrow \mathbb{R} : u \mapsto \frac{1}{2} \int_{\Omega} |\nabla u|^2 + V(x)u^2 - \frac{1}{p} \int_{\Omega} |u|^p.$$

We consider  $\mathcal{N}_p := \{u \neq 0 : \mathcal{E}'_p(u)(u) = 0, \mathcal{E}'_p(u)(v) = 0, \text{ for any } v \in H^-\}$ , where  $H^-$  is the negative spectrum subspace of  $-\Delta + V$ .

- For the  $q$ -Laplacian (i.e.  $\Delta_q u = \text{div}(|\nabla u|^{q-2} \nabla u)$ ) problem

$$\begin{cases} -\Delta_q u(x) = |u(x)|^{p-2}u(x), & x \in \Omega, \\ u(x) = 0, & x \in \partial\Omega, \end{cases}$$

accumulation points, for  $p \rightarrow q$ ,  $u_q$  of **GS** or **LENS** are functions solving  $-\Delta_q u = |u|^{q-2}u$  with  $u = 0$  on  $\partial\Omega$  (\*). If there exist non-zero solutions at (\*), i.e. 1 is eigenvalue of  $\Delta_q$ ,  $u_q$  is different from 0 [9].

- Symmetry results work on nonlinearities  $u|u|^{p-2} + (p-2)u|u|^{q-2}$ ,  $u(e^{u^2} - 1)^{p-2}$  or  $u \left( \sum_{i=1}^k \alpha_i |u|^{\beta_i(p-2)} \right)$ , or for  $\partial u = 0$  on  $\partial\Omega$  [4].

## Future works

- Symmetry breaking curves
- Symmetries, uniqueness, and multiplicity of solutions for Neumann boundary conditions and  $p$  large
- Wrinkling shapes of thin elastic films on water or polymer substrates

